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Daniel Becker Asymptotic Safety & Background Independence

Outline



Asymptotic Safety & Background Independence Single-Metric & Bi-Metric Approximations

Bi-Metric Einstein-Hilbert Truncation: Results Asymptotic Safety & Background Independence The running UV-attractor Single-Metric vs. Bi-Metric: a comparison

3 Conclusion

Asymptotic Safety (UV)

Non-perturbative renormalizability : \Leftrightarrow

• Existenz of a (non-trivial) UV-fixed point of FRGE ...

Evidence for topology-, field-, gauge group-, constraints changes \checkmark

- \ldots with finite dimensional critical hypersurface \mathscr{S}_{UV}

Tests: f(R)-truncations \checkmark

Background Independence

In quantum gravity spacetime is dynamical \Rightarrow Background Independence

In practice two different approaches:

1 No background at all \Rightarrow CDT, LQG, ...

Construct spacetime out of 'Nothing'

 $\ensuremath{ 2 \ } \ensuremath{ \mathsf{Generic}} \ensuremath{ \mathsf{background}} \Rightarrow \mathsf{EAA}$

Choose arbitrary background $\bar{g}_{\mu
u}$ on which dynamical field $g_{\mu
u}$ propagate

Physical sector, $\Gamma_{k=0}$ should be independent on specific $\bar{g}!$

 $\Rightarrow \text{Split-symmetry condition:} \quad \Gamma^{\text{grav}}_{k=0}[g, \bar{g} + \delta \bar{g}] = \Gamma^{\text{grav}}_{k=0}[g, \bar{g}]$

A Global Requirement

Requirements on the FRGE flow for quantum gravity:

- In the UV: Asymptotic Safety (NGFP, $\dim \mathscr{S}_{\text{UV}} < \infty)$
- In the IR: Restoration of split-symmetry $\Gamma_{k=0}$

Question:

Existence of RG trajectories satisfying both conditions (global issue!)

Introduction

Asymptotic Safety & Background Independence

Single-Metric & Bi-Metric Approximations

Functional Renormalization Group Equation

Nonlinear functional differential equation:

$$\partial_t \Gamma_k[g,\bar{g}] = \frac{1}{2} \mathsf{STr} \left\{ \left(\Gamma_k^{(2)}[g,\bar{g}] + R_k[\bar{g}] \right)^{-1} \partial_t R_k[\bar{g}] \right\}$$

Evaluation steps:

- **1** Functional: Reduction to ODE for couplings (LHS & RHS)
- **2** Hessian $\Gamma_k^{(2)}$: w.r.t dynamical fields, $g_{\mu\nu}$, ξ , $\bar{\xi}$ (RHS)
- S Nonlinear: Invert operator (RHS)
- **④ Trace** STr: heat kernel techniques, ... (RHS)

The Bi-Metric Einstein-Hilbert Truncation

1 Effective Average Action $\Gamma_k[g, \bar{g}]$:

(Invariants up to second order in the derivatives)

$$\partial_{t}\Gamma_{k}[g,\bar{g}] = \partial_{t} \int \sqrt{\bar{g}} \left(\Lambda_{k}^{\mathsf{B}} Z_{k}^{\mathsf{B}} + Z_{k}^{\mathsf{B}} \bar{R} + Z_{k}^{a} R \right) + \sqrt{g} \left(\gamma_{k}^{c} + Z_{k}^{c} \bar{R} + Z_{k}^{d} R \right) \\ + \partial_{t} \sum_{n \neq 0} \gamma_{k}^{(n)} \int \left(\frac{\sqrt{\bar{g}}}{\sqrt{g}} \right)^{n} (\cdots) + \partial_{t} \Gamma_{k}^{\mathsf{gf}}[g,\bar{g}] + \partial_{t} \Gamma_{k}^{\mathsf{gh}}[g,\bar{g},\xi,\bar{\xi}]$$

2 Hessian $\Gamma_k^{(2)}[g,\bar{g}]$:

Inversion:

The Bi-Metric Einstein-Hilbert Truncation

1 Effective Average Action $\Gamma_k[g, \bar{g}]$:

(Invariants up to second order in the derivatives)

$$\partial_{t}\Gamma_{k}[g,\bar{g}] = \partial_{t} \int \sqrt{\bar{g}} \left(\Lambda_{k}^{\mathsf{B}} Z_{k}^{\mathsf{B}} + Z_{k}^{\mathsf{B}} \bar{R} + Z_{k}^{a} R \right) + \sqrt{g} \left(\gamma_{k}^{c} + Z_{k}^{c} \bar{R} + Z_{k}^{d} R \right) \\ + \partial_{t} \sum_{n \neq 0} \gamma_{k}^{(n)} \int \left(\frac{\sqrt{\bar{g}}}{\sqrt{g}} \right)^{n} (\cdots) + \partial_{t} \Gamma_{k}^{\mathsf{gf}}[g,\bar{g}] + \partial_{t} \Gamma_{k}^{\mathsf{gh}}[g,\bar{g},\xi,\bar{\xi}]$$

2 Hessian
$$\Gamma_k^{(2)}[g, \bar{g}]$$
:
 $\Gamma_k^{(2)}[g, \bar{g}, \xi, \bar{\xi}] = \begin{pmatrix} \Gamma_{k\,gg}^{(2)} & (\cdots)\bar{\xi} \\ (\cdots)\xi & (\cdots) \end{pmatrix}$

3 Inversion:

4 Trace STr:

The Bi-Metric Einstein-Hilbert Truncation

1 Effective Average Action $\Gamma_k[g, \bar{g}]$:

(Invariants up to second order in the derivatives)

$$\partial_{t}\Gamma_{k}[g,\bar{g}] = \partial_{t} \int \sqrt{\bar{g}} \left(\Lambda_{k}^{\mathsf{B}} Z_{k}^{\mathsf{B}} + Z_{k}^{\mathsf{B}} \bar{R} + Z_{k}^{a} R \right) + \sqrt{g} \left(\gamma_{k}^{c} + Z_{k}^{c} \bar{R} + Z_{k}^{d} R \right) \\ + \partial_{t} \sum_{n \neq 0} \gamma_{k}^{(n)} \int \left(\frac{\sqrt{\bar{g}}}{\sqrt{g}} \right)^{n} (\cdots) + \partial_{t} \Gamma_{k}^{\mathsf{gf}}[g,\bar{g}] + \underline{\partial_{t}} S^{\mathsf{gh}}[g,\bar{g},\bar{\xi},\bar{\xi}]$$

The Bi-Metric Einstein-Hilbert Truncation

1 Effective Average Action $\Gamma_k[g, \bar{g}]$:

(Invariants up to second order in the derivatives)

$$\partial_t \Gamma_k[g, \bar{g}] = \partial_t \int \sqrt{\bar{g}} \left(\Lambda_k^{\mathsf{B}} Z_k^{\mathsf{B}} + Z_k^{\mathsf{B}} \bar{R} + Z_k^{a} R \right) + \sqrt{g} \left(\gamma_k^c + Z_k^c \bar{R} + Z_k^d R \right) \\ + \partial_t \sum_{n \neq 0} \gamma_k^{(n)} \int \left(\frac{\sqrt{\bar{g}}}{\sqrt{g}} \right)^n (\cdots) + \partial_t \Gamma_k^{\mathsf{gf}}[g, \bar{g}]$$

2 Hessian $\Gamma_k^{(2)}[g,\bar{g}]$:

$$\Gamma_{k}^{(2)}[g,\bar{g}] = (\cdots) D_{\mu} D_{\nu} + (\cdots) D_{\mu} \bar{D}_{\nu} + (\cdots) \bar{D}_{\mu} D_{\nu} + (\cdots) \bar{D}_{\mu} \bar{D}_{\nu}$$

The Single-Metric Approximation $\mathsf{FRGE}|_{g=\bar{g}}$

RHS: $\operatorname{STr}\{(\cdots)\}|_{g=\overline{g}} = \operatorname{STr}\{(\cdots)\overline{D}_{\mu}\overline{D}_{\nu}\}$

LHS:

$$\partial_{t}\Gamma_{k}[g=\bar{g},\bar{g}] = \partial_{t}\int\!\!\sqrt{\bar{g}}\left(\Lambda_{k}^{\mathsf{B}}Z_{k}^{\mathsf{B}} + Z_{k}^{\mathsf{B}}\bar{R} + Z_{k}^{a}\bar{R}\right) + \sqrt{\bar{g}}\left(\gamma_{k}^{c} + Z_{k}^{c}\bar{R} + Z_{k}^{d}\bar{R}\right) \\ + \partial_{t}\sum_{n\neq0}\gamma_{k}^{(n)}\int\left(\frac{\sqrt{\bar{g}}}{\sqrt{\bar{g}}}\right)^{n}(\cdots) + \partial_{t}\Gamma_{k}^{\mathsf{gf}}[\bar{g},\bar{g}]$$

The Single-Metric Approximation $\mathsf{FRGE}|_{g=\bar{g}}$

RHS:
$$\operatorname{STr}\{(\cdots)\}|_{g=\overline{g}} = \operatorname{STr}\{(\cdots)\overline{D}_{\mu}\overline{D}_{\nu}\}$$

LHS:

$$\partial_t \Gamma_k[g = \overline{g}, \overline{g}] = \partial_t \int \sqrt{\overline{g}} \underbrace{\left(\underbrace{\Lambda_k^{\mathsf{B}} Z_k^{\mathsf{B}} + \gamma_k^a + \sum_n \gamma_k^{(n)}}_{\equiv \Lambda_k^{\mathsf{sm}} / G_k^{\mathsf{sm}}} + \int \sqrt{\overline{g}} \overline{R} \underbrace{\left(\underbrace{Z_k^{\mathsf{B}} + Z_k^b + Z_k^c + Z_k^d + (\cdots)}_{\equiv 1 / G_k^{\mathsf{sm}}} \right)}_{\equiv 1 / G_k^{\mathsf{sm}}}$$

Huge loss of information:

- \Rightarrow ∞ number of couplings degenerates to only 2, $\Lambda_k^{\rm sm}$ and $G_k^{\rm sm}$
- \Rightarrow Bi-metric character is lost! (split-symmetry enforced for all k)

The Conformal Projection Technique $\mathsf{FRGE}|_{q=e^{2\Omega}\bar{q}}$

RHS:
$$STr\{(\cdots)\}|_{g=e^{2\Omega}\overline{g}} = STr\{(\cdots)\overline{D}_{\mu}\overline{D}_{\nu} + \Omega(\cdots)\overline{D}_{\mu}\overline{D}_{\nu} + \mathcal{O}(\Omega^2)\}$$

LHS:

$$\partial_{t}\Gamma_{k}[g = e^{2\Omega}\bar{g}, \bar{g}] = \partial_{t} \int \sqrt{\bar{g}} \left(\Lambda_{k}^{\mathsf{B}} Z_{k}^{\mathsf{B}} + Z_{k}^{\mathsf{B}} \bar{R} + Z_{k}^{b} e^{-2\Omega} \bar{R}\right) + (\cdots) \\ + \partial_{t} \int e^{d\Omega} \sqrt{\bar{g}} \left(\Lambda_{k}^{c} + Z_{k}^{c} \bar{R} + Z_{k}^{d} e^{-2\Omega} \bar{R}\right) + \underline{\partial_{t}} \Gamma_{k}^{\mathsf{gf}}[e^{2\Omega} \bar{g}, \bar{g}]$$

The Conformal Projection Technique $\mathsf{FRGE}|_{g=e^{2\Omega}\bar{g}}$ RHS: $\mathsf{STr}\{(\cdots)\}|_{g=e^{2\Omega}\bar{g}} = \mathsf{STr}\{(\cdots)\bar{D}_{\mu}\bar{D}_{\nu} + \Omega(\cdots)\bar{D}_{\mu}\bar{D}_{\nu} + \mathcal{O}(\Omega^2)\}$ LHS:

$$\begin{split} \partial_t \Gamma_k [g = e^{2\Omega} \bar{g}, \bar{g}] &= \partial_t \int \!\!\!\sqrt{\bar{g}} \underbrace{\left(\underline{\Lambda_k^{\mathsf{B}} \mathbf{Z}_k^{\mathsf{B}} + \gamma_k^a + \sum_n \gamma_k^{(n)}}_{\mathbf{Z}_k^{\mathsf{D}} = \frac{\mathbf{A}_k^{\mathsf{D}}}{\mathbf{C}_k^{\mathsf{B}}} + \frac{\mathbf{A}_k^{\mathsf{D} \mathsf{p} \mathsf{n}}}{\mathbf{C}_k^{\mathsf{D} \mathsf{p} \mathsf{n}}} + \int \sqrt{\bar{g}} \bar{R} \underbrace{\left(\underline{\mathcal{Z}_k^{\mathsf{B}} + \mathcal{Z}_k^{\mathsf{b}} + \mathcal{Z}_k^{\mathsf{c}} + \mathcal{Z}_k^{\mathsf{d}} + (\cdots) \right)}_{\equiv \frac{1}{\mathcal{C}_k^{\mathsf{D}}} = \frac{1}{\mathcal{C}_k^{\mathsf{B}}} + \frac{1}{\mathcal{C}_k^{\mathsf{D} \mathsf{p} \mathsf{n}}}} \\ &+ \Omega \partial_t \int \!\!\sqrt{\bar{g}} \underbrace{\left(\underline{\gamma_k^{\mathsf{a}} + \sum_n \gamma_k^{(n)}}_{\mathbf{Z}_k^{\mathsf{D} \mathsf{p} \mathsf{n}}} \right)}_{\equiv \mathbf{A}_k^{\mathsf{D} \mathsf{p} \mathsf{n}}} + \int \sqrt{\bar{g}} \bar{R} \underbrace{\left(\underline{\mathcal{Z}_k^{\mathsf{b}} + \mathcal{Z}_k^{\mathsf{c}} + \mathcal{Z}_k^{\mathsf{d}} + (\cdots) \right)}_{\equiv 1/\mathcal{C}_k^{\mathsf{D} \mathsf{p} \mathsf{n}}} \end{split}$$

- Absorb ∞ number of couplings in 4, $\Lambda_k^{\rm Dyn}$, $G_k^{\rm Dyn}$, $\Lambda_k^{(0)/{\rm B}}$, and $G_k^{(0)/{\rm B}}$
- $1/G_k^{\mathsf{B}}, \Lambda_k^{\mathsf{B}} \neq 0$ break split-symmetry ($\mathcal{O}(\Omega^0) \neq \mathcal{O}(\Omega^p) \ p \geq 1$)

$$\begin{split} 32\pi G_k^{\rm sm} \, \Gamma_k^{(2)} [e^{2\Omega} \bar{g}, \bar{g}]_{gg}^{\mu\nu\rho\sigma} &= -\alpha^{-1} \Big[\bar{g}^{\rho\mu} \bar{g}^{\sigma\nu} + (2\varpi^2 - 1) \bar{g}^{\mu\nu} \bar{g}^{\rho\sigma} \Big] \bar{D}^2 \\ &\quad -\alpha^{-1} (1 - 2\varpi) \Big[\bar{g}^{\rho\sigma} \bar{D}^{\mu} \bar{D}^{\nu} + \bar{g}^{\mu\nu} \bar{D}^{\sigma} \bar{D}^{\rho} \Big] \\ &\quad + \Big(\alpha^{-1} - 1 \Big) (\cdots) \bar{D}^{\mu} \bar{D}^{\nu} \\ &\quad + 2 \, \alpha^{-1} \Big[\bar{g}^{\nu\sigma} \bar{R}^{\mu\rho} + \bar{R}^{\rho\mu\nu\sigma} \Big] + \mathcal{U}^{\mu\nu\rho\sigma} [\bar{g}] \end{split}$$

$$\begin{split} 32\pi G_k^{\rm sm} \, \Gamma_k^{(2)} [e^{2\Omega} \bar{g}, \bar{g}]_{gg}^{\mu\nu\rho\sigma} &= -\alpha^{-1} \Big[\bar{g}^{\rho\mu} \bar{g}^{\sigma\nu} + (2\varpi^2 - 1) \bar{g}^{\mu\nu} \bar{g}^{\rho\sigma} \Big] \bar{D}^2 \\ &\quad -\alpha^{-1} (1 - 2\varpi) \Big[\bar{g}^{\rho\sigma} \bar{D}^{\mu} \bar{D}^{\nu} + \bar{g}^{\mu\nu} \bar{D}^{\sigma} \bar{D}^{\rho} \Big] \\ &\quad + \Big(\alpha^{-1} - 1 \Big) (\cdots) \bar{D}^{\mu} \bar{D}^{\nu} \\ &\quad + 2 \, \alpha^{-1} \Big[\bar{g}^{\nu\sigma} \bar{R}^{\mu\rho} + \bar{R}^{\rho\mu\nu\sigma} \Big] + \mathcal{U}^{\mu\nu\rho\sigma} [\bar{g}] \end{split}$$



$$\begin{split} 32\pi G_k^{\mathsf{Dyn}} \, \Gamma_k^{(2)} [e^{2\Omega} \bar{g}, \bar{g}]_{gg}^{\mu\nu\rho\sigma} &= -\alpha^{-1} \Big[\bar{g}^{\rho\mu} \bar{g}^{\sigma\nu} + (2\varpi^2 - 1) \bar{g}^{\mu\nu} \bar{g}^{\rho\sigma} \Big] \bar{D}^2 \\ &\quad -\alpha^{-1} (1 - 2\varpi) \Big[\bar{g}^{\rho\sigma} \bar{D}^{\mu} \bar{D}^{\nu} + \bar{g}^{\mu\nu} \bar{D}^{\sigma} \bar{D}^{\rho} \Big] \\ &\quad + \Big(\alpha^{-1} - e^{(d-6)\Omega} \Big) (\cdots) \bar{D}^{\mu} \bar{D}^{\nu} \\ &\quad + 2\alpha^{-1} \Big[\bar{g}^{\nu\sigma} \bar{R}^{\mu\rho} + \bar{R}^{\rho\mu\nu\sigma} \Big] + e^{(d-6)\Omega} \, \mathcal{U}^{\mu\nu\rho\sigma} [\bar{g}] \end{split}$$

Single-Metric	Bi-Metric I	
$\varpi = 1/2$	arpi = 1/d, $lpha ightarrow 0$	
$\alpha = 1$	$h_{\mu\nu} = h_{\mu\nu}^{TT} + \bar{D}_{\mu}\chi_{\nu} + \cdots$	
$\Rightarrow STr\{(\cdots)\bar{D}^2\}$	TT-decomposition	

$$\begin{split} 32\pi G_k^{\mathsf{Dyn}} \, \Gamma_k^{(2)} [e^{2\Omega} \bar{g}, \bar{g}]_{gg}^{\mu\nu\rho\sigma} &= -\alpha^{-1} \Big[\bar{g}^{\rho\mu} \bar{g}^{\sigma\nu} + (2\varpi^2 - 1) \bar{g}^{\mu\nu} \bar{g}^{\rho\sigma} \Big] \bar{D}^2 \\ &\quad -\alpha^{-1} (1 - 2\varpi) \Big[\bar{g}^{\rho\sigma} \bar{D}^{\mu} \bar{D}^{\nu} + \bar{g}^{\mu\nu} \bar{D}^{\sigma} \bar{D}^{\rho} \Big] \\ &\quad + \Big(\alpha^{-1} - e^{(d-6)\Omega} \Big) (\cdots) \bar{D}^{\mu} \bar{D}^{\nu} \\ &\quad + 2 \alpha^{-1} \Big[\bar{g}^{\nu\sigma} \bar{R}^{\mu\rho} + \bar{R}^{\rho\mu\nu\sigma} \Big] + e^{(d-6)\Omega} \mathcal{U}^{\mu\nu\rho\sigma} [\bar{g}] \end{split}$$

Single-Metric	Bi-Metric I	Bi-Metric II
$\varpi = 1/2$	arpi = 1/d , $lpha ightarrow 0$	$\varpi = 1/2$, $\alpha' = e^{-(d-6)\Omega}$
$\alpha = 1$	$h_{\mu\nu} = h_{\mu\nu}^{TT} + \bar{D}_{\mu}\chi_{\nu} + \cdots$	$\Gamma_k o \Gamma_k + \Gamma_k^{simplify}$
$\Rightarrow STr\{(\cdots)\bar{D}^2\}$	TT-decomposition	$\Rightarrow STr\{(\cdots)\bar{D}^2\}$

Summary: Introduction

A quantum theory of gravity is described by:

Split-symmetry $\stackrel{k=0}{\longleftarrow}$ RG trajectory $\stackrel{k\to\infty}{\longrightarrow}$ NGFP (with dim $\mathscr{S}_{UV} < \infty$)

Evaluation of the FRGE needs approximation techniques:

- Single-metric approximation: $FRGE|_{q=\bar{q}}$ (Background Independence \times)
- Conformal projection: $\mathsf{FRGE}|_{g=e^{2\Omega}\bar{g}}$ (Background Independence (\checkmark))
- New bi-metric method with $\varpi = 1/2$ and $\alpha' = e^{(6-d)\Omega}$:
 - Close analogy to single-metric calculation
 - Simplifies trace evaluation
 - Justification: $\Gamma_k \to \Gamma_k + \Gamma_k^{\text{simplify}}$

Introduction 0000 000000

Asymptotic Safety & Background Independence

Bi-Metric Einstein-Hilbert Truncation

A Bi-Metric Einstein-Hilbert Truncation

$$\Gamma_k^{\rm grav}[g,\bar{g}] = -\frac{1}{16\pi \, G^{\rm B}{}_k} \int \sqrt{\bar{g}} \, \left(\bar{R} - 2\Lambda^{\rm B}{}_k\right) - \frac{1}{16\pi \, G^{\rm Dyn}{}_k} \int \sqrt{g} \, \left(R - 2\Lambda^{\rm Dyn}{}_k\right)$$

- Asymptotic Safety condition
- Split-symmetry, i.e. $\Gamma_k^{\rm grav}[g,\bar{g}] \stackrel{!}{=} \Gamma_k^{\rm grav}[g,\bar{g}+\delta\bar{g}]$

$$\Rightarrow \qquad \frac{1}{G_k^{\mathsf{B}}} \stackrel{!}{=} 0 , \qquad \frac{\Lambda_k^{\mathsf{B}}}{G_k^{\mathsf{B}}} \stackrel{!}{=} 0$$

- Explore Background Independence $k \to 0$
- Test single-metric reliability $\forall k$

Structure of the Beta-Functions

$$\begin{split} \partial_t G_k^{\mathsf{Dyn}} &= \eta^{\mathsf{Dyn}} (G_k^{\mathsf{Dyn}}, \Lambda_k^{\mathsf{Dyn}}) \ G_k^{\mathsf{Dyn}}, \qquad \partial_t \Lambda_k^{\mathsf{Dyn}} = \beta_\Lambda^{\mathsf{Dyn}} (G_k^{\mathsf{Dyn}}, \Lambda_k^{\mathsf{Dyn}}) \\ & \downarrow \\ \partial_t G_k^{\mathsf{B}} &= \eta^{\mathsf{B}} (G_k^{\mathsf{Dyn}}, \Lambda_k^{\mathsf{Dyn}}, G_k^{\mathsf{B}}) \ G_k^{\mathsf{B}} &= \eta^{\mathsf{B}} (G_k^{\mathsf{B}}; \mathbf{k}) \ G_k^{\mathsf{B}} \\ & \downarrow \\ \partial_t \Lambda_k^{\mathsf{B}} &= \beta_\Lambda^{\mathsf{B}} (G_k^{\mathsf{Dyn}}, \Lambda_k^{\mathsf{Dyn}}, G_k^{\mathsf{B}}, \Lambda_k^{\mathsf{B}}) &= \beta_\Lambda^{\mathsf{B}} (G_k^{\mathsf{B}}, \Lambda_k^{\mathsf{B}}; \mathbf{k}) \end{split}$$

For fixed $k \mapsto (G_k^{\mathsf{Dyn}}, \Lambda_k^{\mathsf{Dyn}})$:

- Non-autonomous system: $\beta_g^{\mathsf{B}}(g_k^{\mathsf{B}}; \mathbf{k})$ and $\beta_{\lambda}^{\mathsf{B}}(g_k^{\mathsf{B}}, \lambda_k^{\mathsf{B}}; \mathbf{k})$
- Fixed points: $\beta_g^{\mathsf{B}}(g_{\bullet}^{\mathsf{B}}(k); \mathbf{k}) = 0 = \beta_{\lambda}^{\mathsf{B}}(g_{\bullet}^{\mathsf{B}}(k), \lambda_{\bullet}^{\mathsf{B}}(k); \mathbf{k})$
- $\Rightarrow \left(G^{\mathsf{B}}_{\bullet}(\mathbf{k}), \Lambda^{\mathsf{B}}_{\bullet}(\mathbf{k})\right) \text{ defines running UV-attractor in B-sector}$

Asymptotic Safety Condition UV



UV-condition fulfilled: \exists NGFP with dim (\mathscr{S}_{UV}) = 4

Introduction 0000 000000 Conclusion

Asymptotic Safety & Background Independence

Split-Symmetry in the IR

$$1/G_k^{\mathsf{B}} = 1/G_{k_0}^{\mathsf{B}} - \int_{k_0}^k \mathrm{d}k' \, k^2 B_1^{\mathsf{B}}(k') \stackrel{!}{=} 0$$

For every $(G_k^{\rm Dyn},\Lambda_k^{\rm Dyn})$ trajectory exists one, and only one, full RG trajectory satisfying the 'final condition',

$$1/G_{\mathsf{B}\mathsf{L}}^{\mathsf{B}}(k=0) = \lim_{k \to 0} 1/G_{\bullet}^{\mathsf{B}}(G_{k}^{\mathsf{D}\mathsf{yn}}, \Lambda_{k}^{\mathsf{D}\mathsf{yn}}; k) = 0$$

 \Rightarrow Fulfills global requirement: $\lim_{k\to\infty}\Gamma_k=\Gamma_*^{\rm NGFP}$ and ${\it G}_{\rm B.L}^{\rm B}>0$

$$\left\{\underbrace{\left(G_{k}^{\mathsf{Dyn}},\Lambda_{k}^{\mathsf{Dyn}},G_{k}^{\mathsf{B}},\Lambda_{k}^{\mathsf{B}}\right)}_{\in\mathcal{T}_{\mathsf{trunc:}}:\,\mathsf{4-dim.}}\mid\underbrace{\left(G_{k}^{\mathsf{Dyn}/\mathsf{B}}>0\right)}_{\mathsf{Asym. Safety: 4-dim.}},\underbrace{\left(G_{k}^{\mathsf{B}}=G_{\mathsf{B.I.}}^{\mathsf{B}}(k)\wedge\Lambda_{k}^{\mathsf{B}}=\Lambda_{\mathsf{B.I.}}^{\mathsf{B}}(k)\right)}_{\mathsf{Back. Ind.: 2-dim}}\right\}$$

The Background Sector (here k = 0)



$$k \mapsto \left(G_k^{\mathsf{Dyn}}, \Lambda_k^{\mathsf{Dyn}} \right) \mathsf{fixed} \quad \Rightarrow$$

k-dependent background vector field!

Bi-Metric Einstein-Hilbert Truncation: Results ○○○○○ ○●○○

The running UV-attractor



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The running UV-attractor

The dual role of the running UV-attractor:

- guarantees non-perturbative renormalizability for $k \to \infty$
- guarantees Background Independence for k
 ightarrow 0













Introduction 0000 000000

Single-Metric vs. Bi-Metric: a comparison

Single-Metric vs. Bi-Metric: a comparison

Single-Metric vs. Bi-Metric: a comparison

Test of (intermediate) Split-symmetry

- \forall trajectories \exists only **isolated points** of **intact** split-symmetry
- $\Rightarrow Single and bi-metric results differ significantly for almost all k (especially in the semi-classical regime)$

The miraculous reliability near the NGFP:

Fixed point region *approximately* split-symmetric

Single-metric reliability

IR by construction and UV (\checkmark) In between (semi-classical, ...) ×



Single-Metric vs. Bi-Metric: a comparison

A suitable RG trajectory



Summary

Asymptotic Safety & Background Independence

 A global requirement that needs bi-metric truncations Bi-EH-truncation: trajectories exists satisfying both conditions B.I. ^{IR} running UV-attractor ^{UV}→ A.S.

(Number of free parameter reduces from 4 to 2)

- Split-symmetry broken for 0 < k < ∞
 ⇒ Single-metric trajectories at best qualitatively correct
- UV Miracle: Split-symmetry *approximately* intact near the NGFP ⇒ Single-metric truncation correspondingly reliable

Conclusion

- Background Independence and A. S. simultaneously satisfied
- Generalize truncations to bi-metric ones (single-metric check)

Thank You!