

Correlation functions on a curved background

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Motivation

- the light of the CMB is polarized (Thomson scattering)

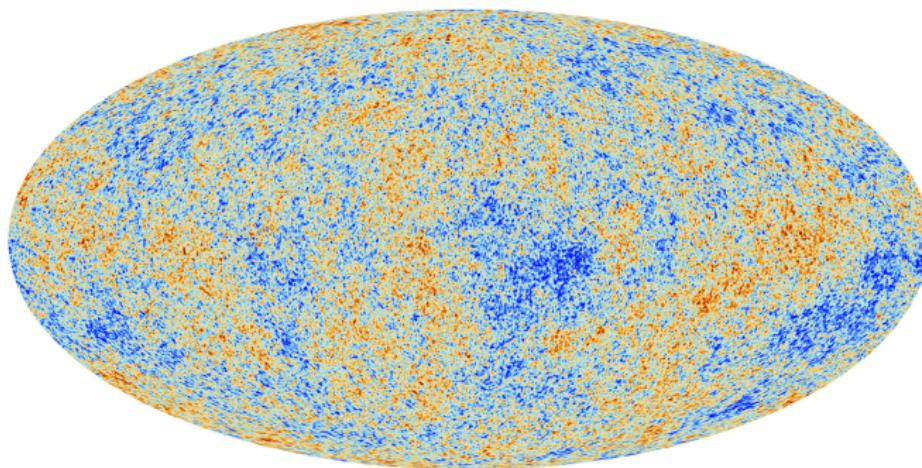


Image: ESA and the Planck Collaboration

Motivation

- the light of the CMB is polarized (Thomson scattering)
- the propagation of this light is influenced by metric perturbations $\delta g_{\mu\nu}$ (scalar and tensor perturbations)
- by measuring the polarization we can infer properties of $\delta g_{\mu\nu}$
- on the other hand, we can calculate these metric perturbations:

$$\delta g_{..}(x) \delta g_{..}(y) \sim \langle h_{..}(x) h_{..}(y) \rangle \sim \Gamma^{(2)}[\bar{g}_{os}] \sim \Gamma_{k=0}^{(0;2)}[0; \bar{g}_{os}]$$

Bi-Metric Nature of Γ_k

Split Symmetry

- linear split: $g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}$
- Einstein-Hilbert action: $S_{\text{EH}}[g] = S_{\text{EH}}[\bar{g} + h]$
- **split symmetry**: $\delta_\varepsilon \bar{g} = -\varepsilon$, $\delta_\varepsilon h = \varepsilon$
 $\Rightarrow \delta_\varepsilon S_{\text{EH}}[g] = 0$
- regulator/gauge-fixing: $\Delta S_k = \Delta S_k[h; \bar{g}]$, $S_{\text{gf}} = S_{\text{gf}}[h; \bar{g}]$
 $\Rightarrow \delta_\varepsilon \Delta S_k[h; \bar{g}] \neq 0$, $\delta_\varepsilon S_{\text{gf}}[h; \bar{g}] \neq 0$

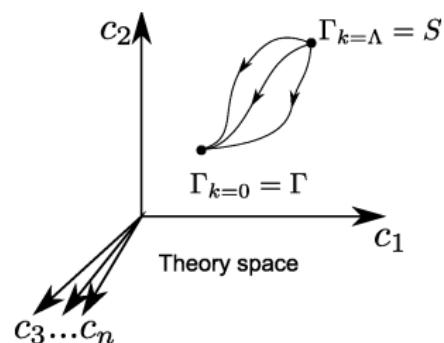
Vertex Expansion

- flow equation:

$$\partial_k \Gamma_k[h; \bar{g}] = \frac{1}{2} S \text{Tr} \left[(\Gamma_k^{(2;0)}[h; \bar{g}] + \mathcal{R}_k[\bar{g}])^{-1} \partial_k \mathcal{R}_k[\bar{g}] \right]$$

$\Rightarrow \partial_k \Gamma_k^{(n;m)}$ obeys hierarchy:

$\Gamma_k^{(2;0)}$	$\Gamma_k^{(3;0)}$	$\Gamma_k^{(4;0)}$	\dots	$\Gamma_k^{(n+2;0)}$
$\Gamma_k^{(2;1)}$	$\Gamma_k^{(3;1)}$	\dots	\dots	$\Gamma_k^{(n+2;1)}$
\vdots	\vdots	\ddots	\ddots	\vdots
$\Gamma_k^{(2;m)}$	$\Gamma_k^{(3;m)}$	\dots	\dots	$\Gamma_k^{(n+2;m)}$



Truncation

- start with Einstein-Hilbert action:

$$S_{\text{EH}}[g] = \frac{1}{16\pi G_N} \int (-R + 2\Lambda)$$

- amend with gauge fixing ($\alpha \rightarrow 0$):

$$S_{\text{gf}}[h; \bar{g}] = \frac{1}{16\pi G_N \alpha} \int \bar{g}^{\mu\nu} F_\mu[h; \bar{g}] F_\nu[h; \bar{g}]$$

$$F_\mu[h; \bar{g}] = \left(\delta_\mu^{(\alpha} \bar{D}^{\beta)} - \frac{1+\beta}{4} \bar{g}^{\alpha\beta} \bar{D}_\mu \right) h_{\alpha\beta}$$

- use them together as generator for vertices:

$$S_{\text{gen}}[h; \bar{g}] = S_{\text{EH}}[\bar{g} + h] + S_{\text{gf}}[h; \bar{g}]$$

- rescale the graviton $h \rightarrow \sqrt{G}h$:

$$S_{\text{EH}}^{(n)}[\bar{g}] \rightarrow G_n^{\frac{n}{2}-1} G_N \left. S_{\text{EH}}^{(n)}[\bar{g}] \right|_{\Lambda \rightarrow \Lambda_n}$$

Truncation

- $\Gamma_k^{(2;0)}:$ $\frac{1}{32\pi} \int h_{\mu\nu} \left(\left(\frac{(\beta-3)^2}{6} \tilde{\Lambda} + \frac{\beta^2}{3} \Lambda_{\text{TL}} \right) \Pi_{\text{Tr}}{}^{\mu\nu}_{\alpha\beta} - \Lambda_{\text{TL}} \Pi_{\text{TL}}{}^{\mu\nu}_{\alpha\beta} \right) h^{\alpha\beta}$
- $\Gamma_k^{(3;0)}:$ $\frac{\sqrt{G_3}}{384\pi} \int h^{\mu\nu} h^{\rho\sigma} \left(\bar{D}_\mu \bar{D}_\nu h_{\rho\sigma} - 14 \bar{D}_{(\mu} \bar{D}_{\rho)} h_{\nu\sigma} - 6 \bar{g}_{\nu\rho} \bar{\Delta} h_{\mu\sigma} \right)$
 $+ \frac{\sqrt{G_3} \Lambda_3}{48\pi} \int h_\mu{}^\nu h_\nu{}^\rho h_\rho{}^\mu$
- $\Gamma_k^{(2;1)}:$ $\frac{1}{32\pi} \int h_{\mu\nu} \left(\frac{1}{3} (1+3\mathcal{R}_{R\text{TL}}) \bar{R} \Pi_{\text{TL}}{}^{\mu\nu}_{\alpha\beta} + (\mathcal{R}_C - 1) \bar{C}^\mu{}^\nu{}_\alpha{}_\beta \right.$
 $+ \mathcal{R}_{S\text{TL}} \Pi_{\text{TL}}{}^{\mu\nu}_{\rho\sigma} \bar{S}_\kappa^\rho \Pi_{\text{TL}}{}^{\sigma\kappa}_{\alpha\beta} + 2 \mathcal{R}_{S\text{Tr}} \bar{g}^{\mu\nu} \bar{S}_{\alpha\beta}$
 $\left. + \mathcal{R}_{\text{Tr}} \bar{R} \Pi_{\text{Tr}}{}^{\mu\nu}_{\alpha\beta} \right) h^{\alpha\beta}$

Regulator

- interplay of Landau limit and degeneracy of scalar sector:

$$\mathcal{P}_{(\sigma h)} = \frac{\bar{\Delta}}{16\alpha} \begin{pmatrix} 12(3-\alpha)\bar{\Delta}^2 & 6(\beta-\alpha)\bar{\Delta} \\ 6(\beta-\alpha)\bar{\Delta} & \beta^2 - 3\alpha \end{pmatrix}$$

$$\mathcal{P}_{(\sigma h)}^{-1} \xrightarrow{\alpha \rightarrow 0} -\frac{4}{3(3-\beta)^2 \bar{\Delta}^3} \begin{pmatrix} \beta^2 & -6\beta\bar{\Delta} \\ -6\beta\bar{\Delta} & 36\bar{\Delta}^2 \end{pmatrix}$$

- in Landau limit propagator is of the form:

$$G_{(\sigma h)} = \begin{pmatrix} \mathcal{O}(\beta^2) & \mathcal{O}(\beta) \\ \mathcal{O}(\beta) & \mathcal{O}(1) \end{pmatrix}$$

- spectrally adjusted regulator:

$$\mathfrak{R}_{(\sigma h)} = \frac{1}{\alpha} \begin{pmatrix} \mathcal{O}(1) & \mathcal{O}(\beta) \\ \mathcal{O}(\beta) & \mathcal{O}(\beta^2) \end{pmatrix}$$

Regulator

- flow equation in Landau limit:

$$\text{Tr}[G_{(\sigma h)} \mathfrak{R}_{(\sigma h)}] = \frac{1}{\alpha} \mathcal{O}(\beta^2) + \mathcal{O}(1)$$

- when everything commutes it is straight forward to ensure that the $\frac{1}{\alpha} \mathcal{O}(\beta^2)$ term vanishes
- but in curved background $[f(\bar{\Delta}), \bar{D}] \neq 0$
 \Rightarrow vanishing of $\frac{1}{\alpha} \mathcal{O}(\beta^2)$ is nontrivial

- well behaved regulator ($\mathfrak{R} \sim S_{\text{gen}}^{(2)} \mathfrak{r}(\frac{\bar{\Delta}}{k^2})$):

$$\begin{aligned} \mathfrak{R} = & \frac{1}{32\pi} \left(-\mathbb{1} \bar{D}^\rho \mathfrak{r}\left(\frac{\bar{\Delta}}{k^2}\right) \bar{D}_\rho + \frac{8\alpha - (1+\beta)^2}{2\alpha} \Pi_{\text{Tr}} \bar{D}^\rho \mathfrak{r}\left(\frac{\bar{\Delta}}{k^2}\right) \bar{D}_\rho \right. \\ & \left. + \frac{1-2\alpha+\beta}{2} \left(\bar{g}^{\cdot\cdot} \bar{D}_{\cdot\cdot} \mathfrak{r}\left(\frac{\bar{\Delta}}{k^2}\right) \bar{D}_{\cdot\cdot} + \bar{D}^{\cdot\cdot} \mathfrak{r}\left(\frac{\bar{\Delta}}{k^2}\right) \bar{D}^{\cdot\cdot} \bar{g}_{\cdot\cdot} \right) + \delta_{\cdot\cdot}^{\cdot\cdot} \bar{D}_{\cdot\cdot} \mathfrak{r}\left(\frac{\bar{\Delta}}{k^2}\right) \bar{D}_{\cdot\cdot} \right) \end{aligned}$$

Results

Fixed Points of the Gauge Parameters

- a direct computation gives:

$$\dot{\alpha} = \mathcal{O}(\alpha^2)$$

$$\dot{\beta} = \mathcal{C}(\alpha, \beta) \cdot \alpha + \mathcal{O}(\alpha^2), \quad \lim_{\alpha \rightarrow 0} |\mathcal{C}(\alpha, \beta)| < \infty$$

- Landau gauge is a fixed point for arbitrary β
- both gauge parameters are exactly marginal couplings

$$\begin{pmatrix} \frac{\partial \dot{\alpha}}{\partial \alpha} & \frac{\partial \dot{\alpha}}{\partial \beta} \\ \frac{\partial \dot{\beta}}{\partial \alpha} & \frac{\partial \dot{\beta}}{\partial \beta} \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ \lim_{\alpha \rightarrow 0} \mathcal{C}(\alpha, \beta) & 0 \end{pmatrix} + \mathcal{O}(\alpha)$$

Fixed Points of the Couplings

- switch to dimensionless variables:

$$g_n = G_n k^2, \quad \lambda_n = \frac{\Lambda_n}{k^2}, \quad \lambda_{\text{TL}} = \frac{\Lambda_{\text{TL}}}{k^2}, \quad \tilde{\lambda} = \frac{\tilde{\Lambda}}{k^2}$$

- note the hierarchy:

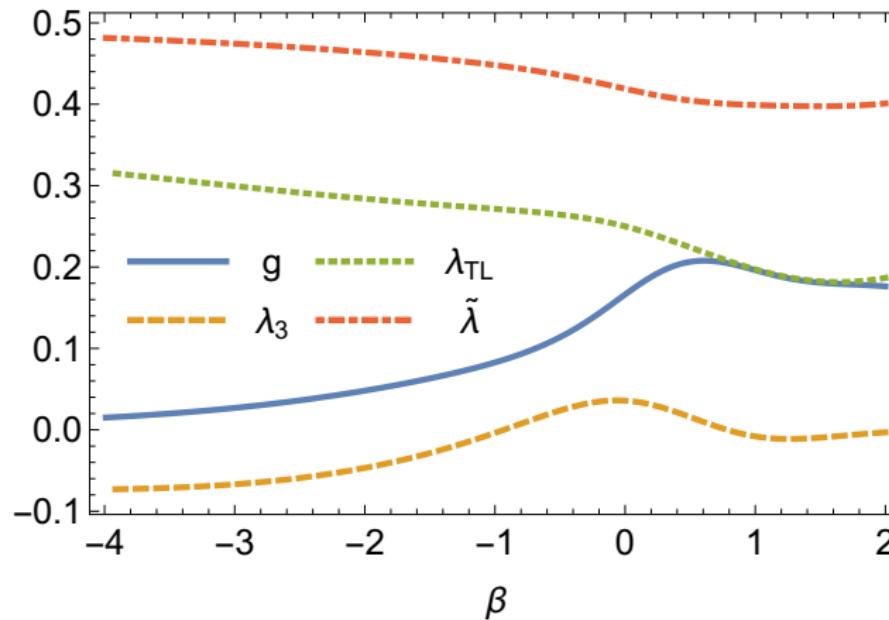
0	$\Gamma_k^{(2;0)}$	$\Gamma_k^{(3;0)}$	$\Gamma_k^{(4;0)}$...	$\Gamma_k^{(n+2;0)}$
1	$\Gamma_k^{(2;1)}$	$\Gamma_k^{(3;1)}$	\dots	\dots	$\Gamma_k^{(n+2;1)}$
\vdots	\vdots	\vdots	\ddots	\ddots	\vdots
m	$\Gamma_k^{(2;m)}$	$\Gamma_k^{(3;m)}$	\dots	\dots	$\Gamma_k^{(n+2;m)}$

- close the flow equation:

$$g_5 = g_4 = g_3 = g, \quad \lambda_5 = \lambda_4 = \lambda_3$$

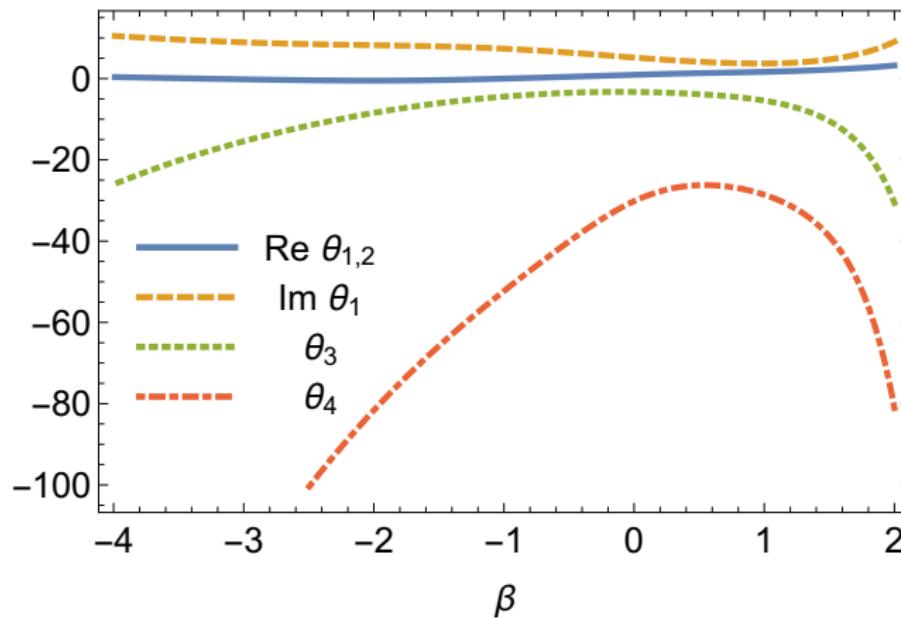
0th Order Curvature Couplings

- $\Gamma_k^{(2;0)}$ and $\Gamma_k^{(3;0)}$



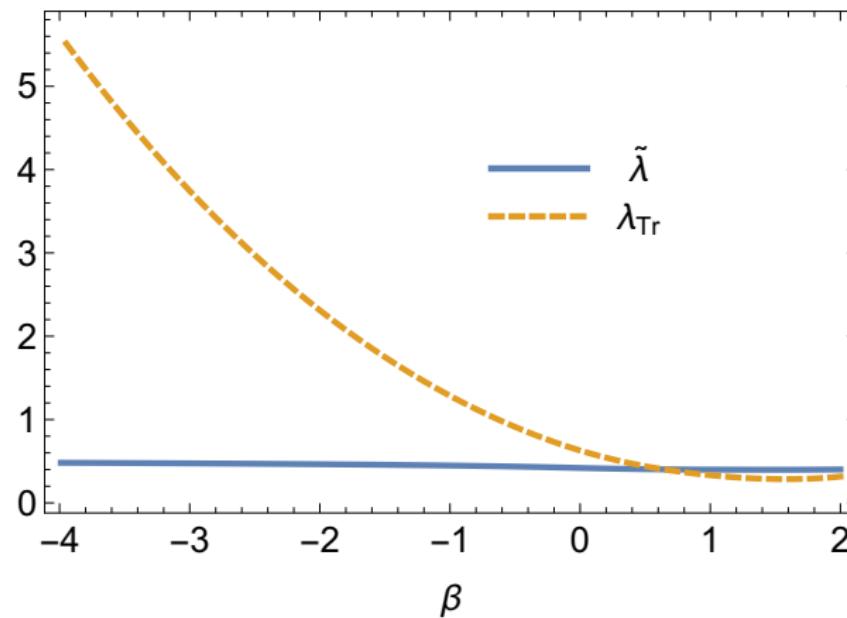
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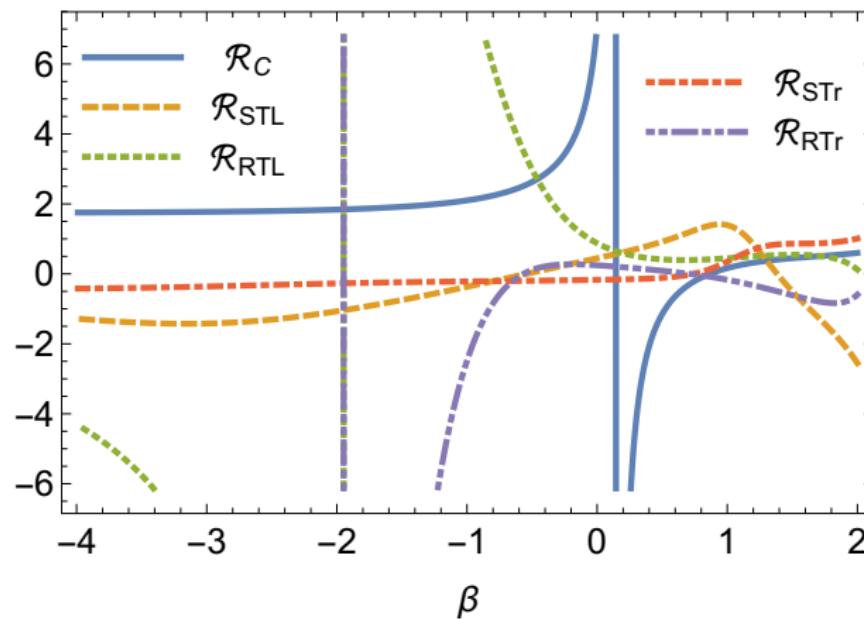
0th Order Curvature Couplings

- $\tilde{\lambda}$ vs. λ_{Tr}



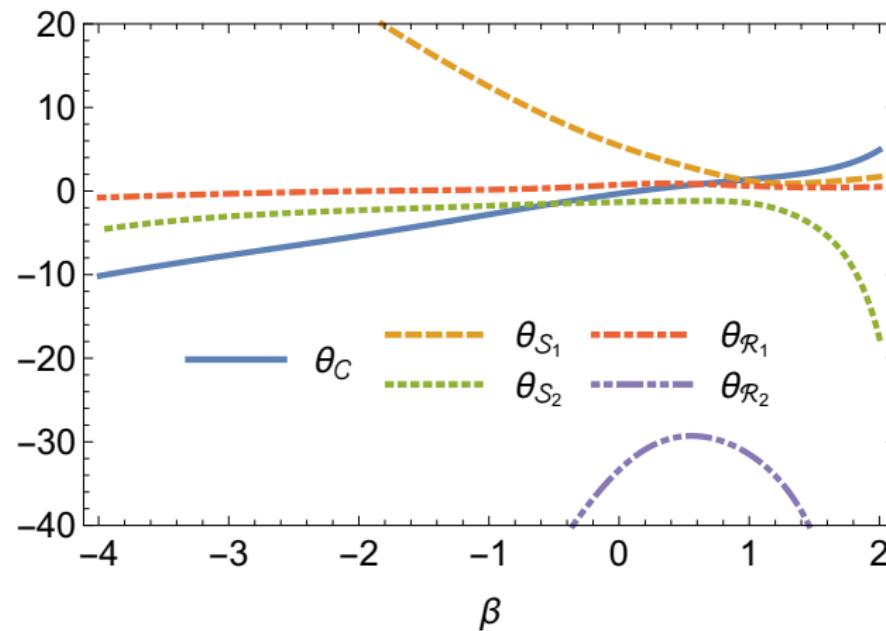
1st Order Curvature Couplings

• $\Gamma_k^{(2;1)}$



1st Order Curvature Couplings

• $\Gamma_k^{(2;1)}$



Summary & Outlook

- existence of unique UV fixed point persists under inclusion of background curvature
- Landau gauge implies fixed point also for β
- regulator choice in curved background
- cosmological post-/predictions
- calculations on self-consistent backgrounds
- inclusion of matter

Thank you for your attention!