

The path integral of unimodular gravity

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Motivation

- ▶ JHEP **1606**, 115 (2016) (N. Ohta, R. Percacci, A. D. Pereira)

$$\int \mathcal{D}\gamma_{\mu\nu} e^{-S[g(\gamma)]}, \quad \int \mathcal{D}\gamma^{\mu\nu} e^{-S[g(\gamma)]}$$
$$\gamma_{\mu\nu} = g_{\mu\nu} (\sqrt{\det g_{\mu\nu}})^{-\frac{m}{1+dm}}, \quad \gamma^{\mu\nu} = g_{\mu\nu} (\sqrt{\det g_{\mu\nu}})^{\frac{m}{1+dm}}$$
$$g_{\mu\nu} = \gamma_{\mu\nu} (\det \gamma_{\mu\nu})^m, \quad g^{\mu\nu} = \gamma^{\mu\nu} (\det \gamma_{\mu\nu})^{-m}$$

- ▶ Not invertible for $m = -1/d$. Unimodular point: $\det g_{\mu\nu} = 1$.
- ▶ The divergences are invariant under a \mathbb{Z}_2 "duality" transformation between $\gamma_{\mu\nu}$ and $\gamma^{\mu\nu}$.

Motivation

- ▶ Unimodular Gravity (UG) : Cosmological constant as an integration constant.
- ▶ **QUG** $\stackrel{?}{=} \mathbf{QGR}$ (Gravity only)
- ▶ Claim: **QUG** = **QGR** with UG as a *S*Diff theory.
(like $O(n)$ and $SO(n)$)

1 loop GR review

► $\hbar = 0$

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G_N^{(d)} T_{\mu\nu}$$

$$S_{Diff} = \frac{1}{16\pi G_N^{(d)}} \int d^d x \sqrt{-g} (R - 2\Lambda) + S_m$$

$$T_{\mu\nu} = -\frac{2}{\sqrt{-g}} \frac{\delta S_m}{\delta g^{\mu\nu}}$$

► $\hbar = 1$

$$Z_{Diff} = \int \mathcal{D}h_{\mu\nu} e^{-S_{Diff}(\bar{g}_{\mu\nu} + h_{\mu\nu})}$$

$$\delta h_{\mu\nu} = \bar{\nabla}_\mu \epsilon_\nu + \bar{\nabla}_\nu \epsilon_\mu$$

$$\bar{R}_{\mu\nu} = \frac{1}{d} \bar{R} \bar{g}_{\mu\nu} \quad (\Lambda = 0)$$

1 loop GR review

- Decomposition:

$$h_{\mu\nu} = h_{\mu\nu}^{\text{TT}} + \bar{\nabla}_\mu \xi_\nu + \bar{\nabla}_\nu \xi_\mu + \left(\bar{\nabla}_\mu \bar{\nabla}_\nu - \frac{1}{d} \bar{g}_{\mu\nu} \bar{\nabla}^2 \right) \sigma + \frac{1}{d} \bar{g}_{\mu\nu} h$$
$$\epsilon_\mu = \epsilon_\mu^{\text{T}} + \bar{\nabla}_\mu \phi$$

- Transformations:

$$\delta h_{\mu\nu}^{\text{TT}} = 0, \quad \delta \xi_\mu = \epsilon_\mu^{\text{T}}, \quad \delta \sigma = 2\phi, \quad \delta h = 2\nabla^2 \phi$$

- One obtains:

$$Z_{\text{Diff}} = \left(\int \mathcal{D}\xi_\mu^{\text{T}} \mathcal{D}\sigma \text{Det}^{1/2} \Delta_{L_0} \right) e^{-S_0} \frac{\text{Det}^{1/2}(\Delta_{L_1} - \frac{2\bar{R}}{d})}{\text{Det}^{1/2}(\Delta_{L_2} - \frac{2\bar{R}}{d})}$$

1 loop GR review

- ▶ The volume of *Diff*'s:

$$V(\text{Diff}) = \int \mathcal{D}\epsilon_\mu = \int \mathcal{D}\epsilon_\mu^T \mathcal{D}\phi \text{Det}^{1/2} \Delta_{L_0}$$

- ▶ Identification:

$$\xi_\mu \leftrightarrow \epsilon_\mu^T, \quad \sigma \leftrightarrow \phi$$

- ▶ One obtains:

$$Z_{\text{Diff}} = V(\text{Diff}) e^{-S_0} \frac{\text{Det}^{1/2}(\Delta_{L_1} - \frac{2\bar{R}}{d})}{\text{Det}^{1/2}(\Delta_{L_2} - \frac{2\bar{R}}{d})}$$

SDiff's

- ▶ Definition:

$$SDiff : \nabla_{\mu}\epsilon^{\mu} = 0, \quad (\bar{\nabla}^2\phi = 0)$$

- ▶ The volume of *SDiff*'s:

$$V(SDiff) = \int \mathcal{D}\epsilon_{\mu}\delta(\nabla_{\mu}\epsilon^{\mu}) = \int \mathcal{D}\epsilon_{\mu}^T \text{Det}^{-1/2}\Delta_{L_0}$$

- ▶ Identification:

$$h \leftrightarrow \bar{\nabla}^2\phi \quad \Rightarrow \quad \boxed{V(Diff) = V(SDiff) \int \mathcal{D}h}$$

- ▶ $V(Q) = \int \mathcal{D}h$, $Q = Diff/SDiff =$ space of volume forms.

QUG as a $SDiff$ theory

- ▶ By definition:

$$\delta_{SDiff} \det g = 2 \det g \nabla_{\mu} \epsilon^{\mu} = 0 \quad (g^{\mu\nu} \delta g_{\mu\nu} = 0)$$

- ▶ The volume form:

$$\text{Vol}_d = \underbrace{\sqrt{-\det g}} \quad dx^0 \wedge dx^1 \wedge \dots \wedge dx^{d-1}$$

Covariant under $Diff'$ s.
Invariant under $SDiff'$ s

QUG as a $SDiff$ theory

- ▶ Exponential Parametrization:

$$g_{\mu\nu} = \bar{g}_{\mu\rho} (e^h)^\rho{}_\nu$$

- ▶ The volume form:

$$\text{Vol}_d = \sqrt{-\det \bar{g}} e^{\text{tr} h} dx^0 \wedge dx^1 \wedge \dots \wedge dx^{d-1}$$

- ▶ We choose: $h_{\mu\nu} = h_{\mu\nu}^T$.
- ▶ This is consistent with the identification $h \leftrightarrow \bar{\nabla}^2 \phi$. Recall that for $SDiff$'s $\bar{\nabla}^2 \phi = 0$.

QUG as a $SDiff$ theory

- ▶ We define the action (ω is fixed):

$$S_{UG} = S_{SDiff} = \frac{1}{16\pi G_N^{(d)}} \int d^d x \omega R + S_m$$

- ▶ $\hbar = 1$

$$\begin{aligned} Z_{SDiff} &= \left(\int \mathcal{D}\xi_\mu^T \right) e^{-S_0} \frac{\text{Det}^{1/2}(\Delta_{L_1} - \frac{2\bar{R}}{d})}{\text{Det}^{1/2}(\Delta_{L_2} - \frac{2\bar{R}}{d})} \frac{1}{\text{Det}^{1/2} \Delta_{L_0}} \\ &= V(SDiff) e^{-S_0} \frac{\text{Det}^{1/2}(\Delta_{L_1} - \frac{2\bar{R}}{d})}{\text{Det}^{1/2}(\Delta_{L_2} - \frac{2\bar{R}}{d})} \end{aligned}$$

- ▶ $Z_{Diff} = Z_{SDiff}$

Hamiltonian analysis of UG

- ▶ $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$
- ▶ UG Lagrangian = Fierz-Pauli Lagrangian with $h = 0$

$$\mathcal{L}_{UG} = -\frac{1}{2}\partial_\alpha h_{\mu\nu}\partial^\alpha h^{\mu\nu} + \partial_\alpha h_\mu{}^\alpha\partial_\beta h^{\mu\beta} - \partial_\alpha h_\mu{}^\alpha\partial^\mu h + \frac{1}{2}\partial_\alpha h\partial^\alpha h$$

- ▶ Decomposition of $h_{\mu\nu}, \epsilon_\mu$

$$h_{00} = -2\phi$$

$$h_{0i} = v_i = v_i^T + \partial_i v$$

$$h_{ij} = h_{ij}^{TT} + \partial_i\zeta_j + \partial_j\zeta_i + \left(\partial_i\partial_j - \frac{1}{d-1}\delta_{ij}\partial^2\right)\tau + \frac{1}{d-1}\delta_{ij}t$$

$$\epsilon_\mu = (\epsilon_0, \epsilon_i^T + \partial_i\epsilon)$$

- ▶ $h = 0$ implies that $2\phi + t = 0$.

Hamiltonian analysis of UG

► Dirac-Bergmann algorithm

GR	UG
d Primary Constraints	$d - 1$ Primary Constraints
d Secondary Constraints	$d - 1$ Secondary Constraints
0 Tertiary Constraints	1 Tertiary Constraint
<hr/>	
$2d$ First class constraints (FCC)	$2d - 1$ First class constraints

► D.o.f in $d = 4$

	GR	UG
(q, p)	2×10	2×9
FCC	8	7
D.o.f = $(q, p) - 2$ FCC	$2+2$	$2+2$

- ▶ $\hbar = 1$

$$Z = \int \mathcal{D}q's \mathcal{D}p's \det M \delta(\phi's) \delta(\chi's) e^{-S[q's, p's]}$$

$$M_{ab} = \{\phi_a, \chi_b\}, \quad \phi's = \text{FCC}, \quad \chi's = \text{Gauge conditions}$$

- ▶ We have

$$\det M_{GR} = (\det \square)^2, \quad \det M_{UG} = (\det \square)^3$$

$$\phi_{GR} \neq \phi_{UG}$$

- ▶ Nevertheless: $Z_{Diff} = Z_{SDiff} = (\det \square_{TT})^{-1/2}$.

Classical UG

$$\delta S_{UG} = \frac{1}{16\pi G_N^{(d)}} \int d^d x \omega (R_{\mu\nu} - 8\pi G^{(d)} \theta_{\mu\nu}) \delta g^{\mu\nu}, \quad (g_{\mu\nu} \delta g^{\mu\nu} = 0)$$

$$R_{\mu\nu} - 8\pi G^{(d)} \theta_{\mu\nu} = \psi g_{\mu\nu}, \quad \theta_{\mu\nu} = \frac{-2}{\omega} \frac{\delta S_m}{\delta g^{\mu\nu}}$$

$$\psi = \tilde{\psi} + 8\pi G^{(d)} f, \quad T_{\mu\nu} = \theta_{\mu\nu} + f g_{\mu\nu} \quad \text{such that} \quad \nabla_\mu T^{\mu\nu} = 0$$

$$R_{\mu\nu} - 8\pi G^{(d)} T_{\mu\nu} = \tilde{\psi} g_{\mu\nu} \quad \Rightarrow \quad \boxed{R_{\mu\nu} - \frac{1}{d} g_{\mu\nu} R = 8\pi G^{(d)} (T_{\mu\nu} - \frac{1}{d} g_{\mu\nu} T)}$$

Classical UG

$$R_{\mu\nu} - \frac{1}{d}g_{\mu\nu}R = 8\pi G^{(d)}\left(T_{\mu\nu} - \frac{1}{d}g_{\mu\nu}T\right)$$

$$\nabla^\mu R_{\mu\nu} - \frac{1}{d}g_{\mu\nu}\nabla^\mu R = -\frac{8\pi G^{(d)}}{d}g_{\mu\nu}\nabla^\mu T, \quad \nabla^\mu G_{\mu\nu} = 0$$

$$\nabla_\mu \left(\frac{d-2}{2d}R + 8\pi G^{(d)}T \right) = 0$$

Then

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G_N^{(d)}T_{\mu\nu}$$

Λ as an integration constant.

UG as a Weyl invariant theory

$$S = \frac{1}{16\pi G_N^{(d)}} \int d^d x \sqrt{-g} R$$

$$g_{\mu\nu} = \gamma_{\mu\nu} \left(\frac{|\gamma|}{\omega^2} \right)^{-1/d}$$

$$S_{UG} = Z_N \int d^d x |\gamma|^{1/d} \omega^{\frac{d-2}{d}} \left[R[\gamma] + \frac{(d-1)(d-2)}{4d^2} (|\gamma|^{-1} \nabla |\gamma| - 2\omega^{-1} \nabla \omega)^2 \right]$$

UG and the Cosmological Constant Problem

$$\text{GR} = \begin{cases} G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G_N^{(d)} T_{\mu\nu} \\ S_{\text{Diff}} = \frac{1}{16\pi G_N^{(d)}} \int d^d x \sqrt{-g} (R - 2\Lambda) + S_m \end{cases}$$

$$\text{UG} = \begin{cases} G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G_N^{(d)} T_{\mu\nu} \\ S_{\text{Diff}} = \frac{1}{16\pi G_N^{(d)}} \int d^d x \omega R + S_m \end{cases}$$

$$S_{\text{Diff}} \supset - \int d^d x \sqrt{-g} (2\Lambda + V(\phi))$$

$$S_{\text{SDiff}} \supset - \int d^d x \omega V(\phi)$$

Thank you