

Higgs scalar potential in asymptotically safe quantum gravity

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Phys.Rev. D99 (2019) no.8, 086010



Alexander von Humboldt
Stiftung / Foundation

Plan

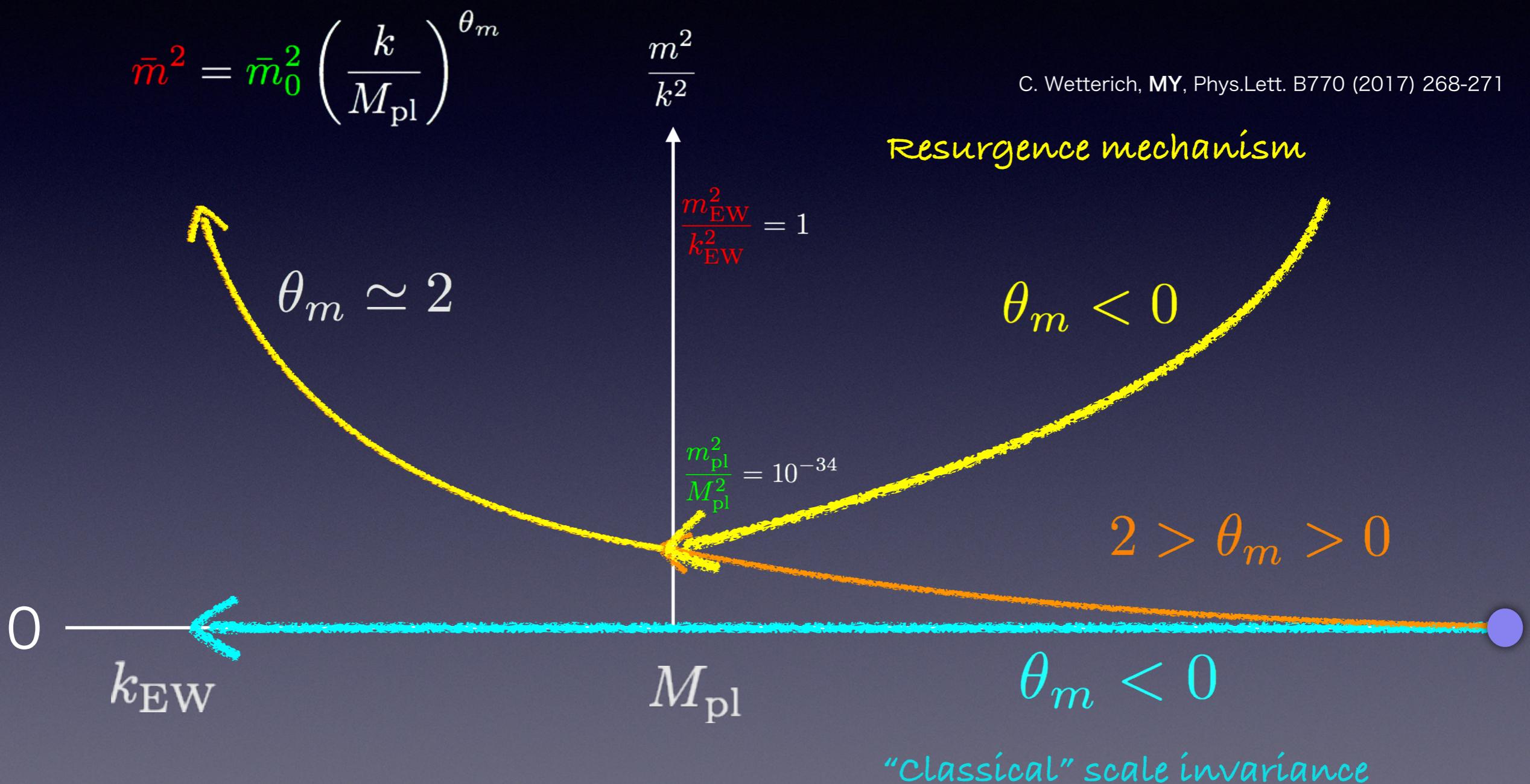
- Quantum scaling
- Gauge invariant flow equation
- Quantum gravity effects on Higgs scalar potential

Quantum Scaling

- Due to graviton fluctuation, a large anomalous dimension could be induced above the Planck scale.
- Energy Scalings deviate from canonical ones.
- Critical exponent (effective dimension) of a scalar mass term

$$\theta_m = 2 - \gamma_{\text{SM}} - \frac{5g^*}{3\pi(1-v^*)^2} \quad (0)$$

RG flow of scalar mass



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Gauge invariant flow equation

- The gauge invariant flow equation

$$\partial_t \Gamma_k = \pi_k - \delta_k \quad (1)$$

- π_k : Physical fluctuations $\pi_k = \frac{1}{2} \text{Str} \{ \partial_t R_P G_P \}$
$$(\bar{\Gamma}_k^{(2)} + R_P) G_P = P^T$$

- δ_k : a universal measure contribution
- See “C. Wetterich Nucl.Phys. B931 (2018) 262-282” for his derivations.

In background formalism

- The metric $g_{\mu\nu} = \delta_{\mu\nu} + h_{\mu\nu}$

- The metric fluctuation

Clear relations with the Bardeen potentials.

C.Wetterich, Phys.Rev. D95 (2017) no.12, 123525

$$h_{\mu\nu} = Ph_{\mu\nu} + (1 - P)h_{\mu\nu}$$

$$= f_{\mu\nu} + a_{\mu\nu}$$

Physical metric fluctuation

$$q^\mu f_{\mu\nu} = 0$$

Gauge metric fluctuation

$$f_{\mu\nu} = t_{\mu\nu} + \frac{1}{3} \left(\delta_{\mu\nu} - \frac{q_\mu q_\nu}{q^2} \right) \sigma$$

$$a_{\mu\nu} = i (q_\mu \kappa_\nu + q_\nu \kappa_\mu) + \frac{q_\mu q_\nu}{q^2} u$$

In background formalism

- Einstein-Hilbert truncation + scalar theory

$$\Gamma_k = -\frac{M_p^2}{2} \int d^4x \sqrt{g} R + \int d^4x \sqrt{g} \left[U(\rho) + \frac{Z_\phi}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi \right] + S_{\text{gf}} + S_{\text{gh}} \quad (2)$$

$$S_{\text{gf}} = \frac{1}{2\alpha} \int d^4x \sqrt{g} g^{\mu\nu} \Sigma_\mu \Sigma_\nu \quad \Sigma_\mu = \bar{\nabla}^\nu h_{\nu\mu} - \frac{\beta+1}{4} \bar{\nabla}_\mu h$$

$$S_{\text{gh}} = - \int d^4x \sqrt{g} C_\mu \left[\bar{g}^{\mu\rho} \bar{\nabla}^2 + \frac{1-\beta}{2} \bar{\nabla}^\mu \bar{\nabla}^\rho + \bar{R}^{\mu\rho} \right] C_\rho$$

- Evaluate the Hessians for

$$h_{\mu\nu} = t_{\mu\nu} + \frac{1}{3} \left(\delta_{\mu\nu} - \frac{q_\mu q_\nu}{q^2} \right) \sigma + i (q_\mu \kappa_\nu + q_\nu \kappa_\mu) + \frac{q_\mu q_\nu}{q^2} u$$

$$C_\mu = C_\mu^\perp + iq_\mu C$$
$$\bar{C}_\mu = \bar{C}_\mu^\perp + iq_\mu \bar{C}$$

Hessian for spin 0 modes

$$\Gamma_{(00)}^{(2)} = \begin{matrix} \sigma & u & \phi \\ \sigma & \left(\Gamma_{(00)}^{(2)} \right)_{\text{grav}} & \frac{U' \phi}{2} \\ \phi & \frac{U' \phi}{2} & \frac{U' \phi}{2} \\ \end{matrix} \quad (3)$$

$$\left(\Gamma_{(00)}^{(2)} \right)_{\text{grav}} = \begin{matrix} \sigma & u \\ \sigma & -\frac{M_p^2}{6} \left(q^2 - \frac{U}{2M_p^2} \right) + \frac{(\beta+1)^2}{16\alpha} q^2 & \frac{U}{4} + \frac{(\beta+1)(\beta-3)}{16\alpha} q^2 \\ u & \frac{U}{4} + \frac{(\beta+1)(\beta-3)}{16\alpha} q^2 & -\frac{U}{4} + \frac{(\beta-3)^2}{16\alpha} q^2 \end{matrix} \quad (4)$$

$$\frac{1}{2} \text{Tr} \frac{\partial_t \mathcal{R}_k}{\Gamma_{(00)}^{(2)} + \mathcal{R}_k} \xrightarrow[\alpha \rightarrow 0]{\beta = -1} \frac{1}{2} \text{Tr} \frac{\partial_t \mathcal{R}_k}{\Gamma_{(\text{ph})}^{(2)} + \mathcal{R}_k} + \frac{1}{2} \text{Tr} \frac{\partial_t \mathcal{R}_k}{\Gamma_{(\text{gauge})}^{(2)} + \mathcal{R}_k}$$

$$\left(\Gamma_{(00)}^{(2)} \right)_{\text{ph}} = \begin{matrix} \sigma & \phi \\ \sigma & -\frac{M_p^2}{6} \left(q^2 - \frac{U}{2M_p^2} \right) & U' \phi / 2 \\ \phi & U' \phi / 2 & Z_\phi q^2 + U' + 2\rho U'' \end{matrix} \quad (5)$$

Each contribution

- Hessians

Spin 2 (graviton)

$$\left(\Gamma_{(tt)}^{(2)}\right)^{\mu\nu\rho\sigma} = \frac{M_p^2}{4} \left[q^2 - \frac{2U}{M_p^2} \right] P^{(t)\mu\nu\rho\sigma}$$

Spin 0

$$\left(\Gamma_{(00)}^{(2)}\right)_{\text{ph}} = \begin{pmatrix} -\frac{M_p^2}{6} \left(q^2 - \frac{U}{2M_p^2} \right) & U'\phi/2 \\ U'\phi/2 & Z_\phi q^2 + U' + 2\rho U'' \end{pmatrix}$$

Physical contributions

Spin 1

$$\left(\Gamma_{(\kappa\kappa)}^{(2)}\right)^{\mu\nu} = q^4 P^{(v)\mu\nu}$$

$$J_{\text{spin } 1}$$

$$\left(\Gamma_{(\bar{C}^\perp C^\perp)}^{(2)}\right)^{\mu\nu} = q^2 P^{(v)\mu\nu}$$

Spin 0

$$\Gamma_{(uu)}^{(2)} = q^2$$

$$\Gamma_{(\bar{C}C)}^{(2)} = q^4$$

$$J_{\text{spin } 0}$$

- RG flow

$$\pi_2 = \frac{1}{2} \text{Tr}_{(2)} \left. \frac{\partial_t \mathcal{R}_k}{\Gamma_k^{(2)} + \mathcal{R}_k} \right|_{tt} = \frac{5}{2} \int_q \frac{\partial_t (M_p^2 R_k)}{M_p^2 (P_k - k^2 v)}$$

$$\pi_0 = \frac{1}{2} \int_q \tilde{\partial}_t \ln \left\{ 3\rho U'^2 + \left(M_p^2 P_k - \frac{U}{2} \right) (Z_\phi P_k + U' + 2\rho U'') \right\}$$

$$\delta_k^{(1)} = \frac{1}{2} \text{Tr}_{(1)} \left. \frac{\partial_t \mathcal{R}_k}{\Gamma_k^{(2)} + \mathcal{R}_k} \right|_{\kappa\kappa} + J_{\text{spin } 1} = \frac{1}{2} \text{Tr}_{(1)} \frac{\partial_t P_k}{P_k}$$

$$-\epsilon_k^{(1)} = -\text{Tr}_{(1)} \left. \frac{\partial_t \mathcal{R}_k}{\Gamma_k^{(2)} + \mathcal{R}_k} \right|_{\bar{C}^\perp C^\perp} = -\text{Tr}_{(1)} \frac{\partial_t P_k}{P_k}$$

$$\epsilon_k^{(1)} = 2\delta_k^{(1)}$$

Gauge and ghost contributions

$$\delta_k^{(0)} = \frac{1}{2} \text{Tr}_{(0)} \left. \frac{\partial_t \mathcal{R}_k}{\Gamma_k^{(2)} + \mathcal{R}_k} \right|_{(uu)} = \frac{1}{2} \text{Tr}_{(0)} \frac{\partial_t P_k}{P_k}$$

$$10^{-\epsilon_k^{(0)}} = -\text{Tr}_{(0)} \left. \frac{\partial_t \mathcal{R}_k}{\Gamma_k^{(2)} + \mathcal{R}_k} \right|_{\bar{C}C} + J_{\text{spin } 0} = -\text{Tr}_{(0)} \frac{\partial_t P_k}{P_k}$$

$$\epsilon_k^{(0)} = 2\delta_k^{(0)}$$

RG flow

- To summarise...

$$\begin{aligned}\partial_t \Gamma_k &= \pi_2 + \pi_0 + (\delta_k^{(1)} - \epsilon_k^{(1)}) + (\delta_k^{(0)} - \epsilon_k^{(0)}) \\ &= \pi_2 + \pi_0 - \delta_k^{(1)} - \delta_k^{(0)}\end{aligned}$$

$\epsilon_k^{(1)} = 2\delta_k^{(1)}$
 $\epsilon_k^{(0)} = 2\delta_k^{(0)}$

$$\text{DOF} \quad 5 + 2 - 3 - 1 = 2 + 1$$

↑
TT mode *σ and φ modes*

Two graviton helicities+scalar (ϕ)

Note

- York decomposition

$$h_{\mu\nu} = \textcolor{red}{t}_{\mu\nu} + i(q_\mu \kappa_\nu + q_\nu \kappa_\mu) - \left(q_\mu q_\nu - \frac{1}{4} \eta_{\mu\nu} q^2 \right) s + \frac{1}{4} \eta_{\mu\nu} \textcolor{green}{h}$$

- The gauge choice $\beta = 0, \alpha \rightarrow 0$

$$\partial_t \Gamma_k = \textcolor{red}{\pi}_2 + \textcolor{green}{\pi}_0 - \delta_k^{(1)} - \delta_k^{(0)}$$

TT mode $\text{Trace mode } h = \delta^{\mu\nu} h_{\mu\nu}$

- The trace mode does not satisfy the transverse condition.

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RG flow for effective scalar potential

SO(N) scalar and free particles

$$\begin{aligned}
 \partial_t \tilde{U} = & -4\tilde{U} + (2 + \eta_\phi)\tilde{\rho}\tilde{U}' \\
 & + \frac{1}{24\pi^2} \left(1 - \frac{\eta_g}{8}\right) \left[\frac{5}{1-v} + \frac{1}{1-v/4} \right] + \frac{\Delta N - 4}{32\pi^2} \\
 & + \frac{1}{32\pi^2} \left(1 - \frac{\eta_\phi}{6}\right) \left[\frac{1}{1+\tilde{U}'+2\tilde{\rho}\tilde{U}''} + \frac{N-1}{1+\tilde{U}'} \right]
 \end{aligned} \tag{6}$$

graviton physical mode
Spin 0

Measure contribution


$$v(\rho) = \frac{2U(\rho)}{M_p^2 k^2} = \frac{2\tilde{U}(\rho)}{\tilde{M}_p^2}$$

$$\Delta N = \Delta N_S + 2N_V - 2N_F$$

Critical exponent

- The effective potential

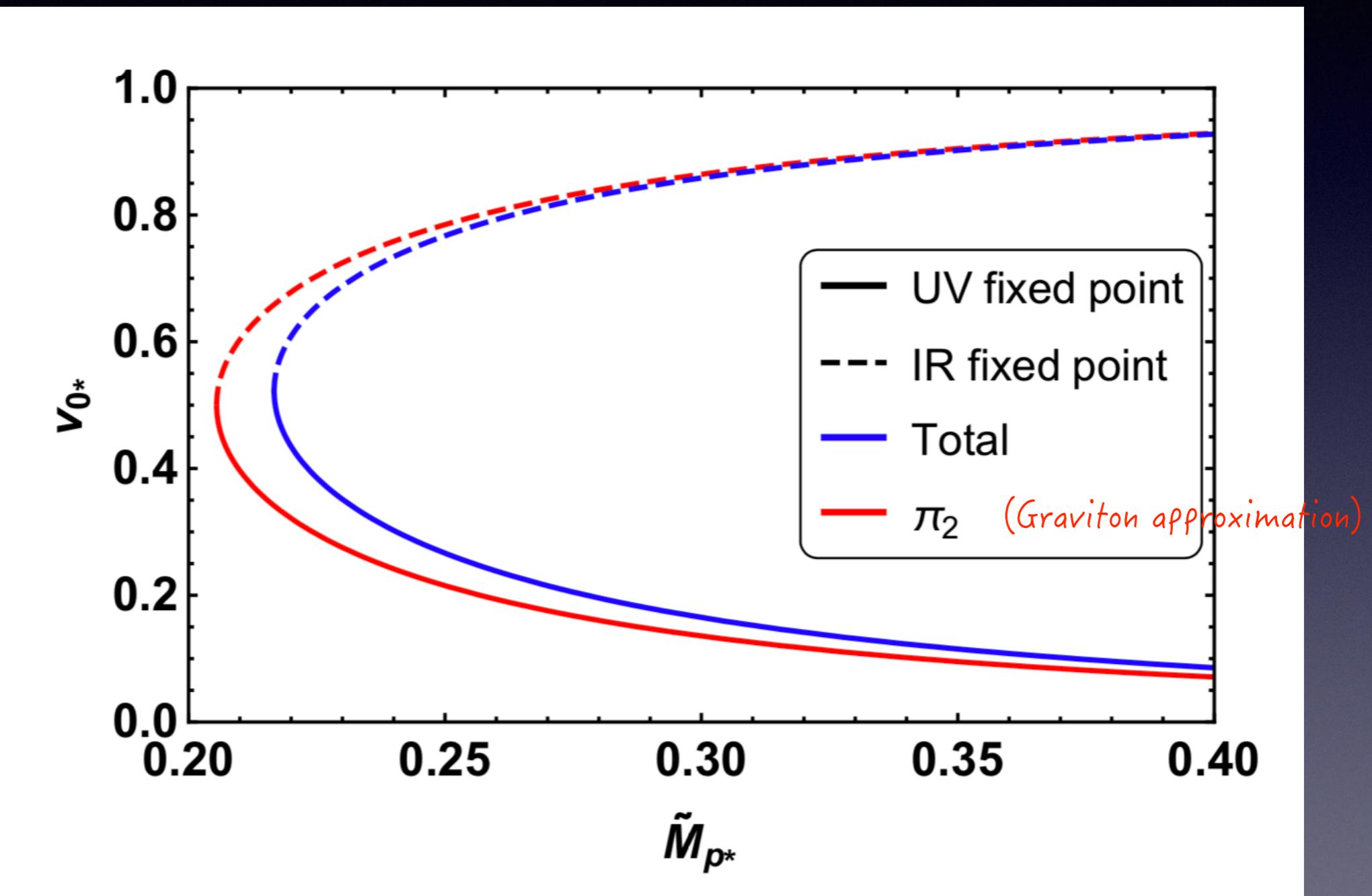
$$\tilde{U} = \tilde{V} + \tilde{m}_H^2 \tilde{\rho} + \frac{\tilde{\lambda}_H}{2} \tilde{\rho}^2$$

- Stability matrix at a Gaussian-matter FP

$$T = \begin{pmatrix} 4 - A & \frac{1}{4\pi^2 \tilde{M}_p^2} & 0 \\ 0 & 2 - A & \frac{3}{16\pi^2} \\ 0 & 0 & -A \end{pmatrix} \quad A = \frac{1}{48\pi^2 \tilde{M}_p^2} \left[\frac{20}{(1 - v_0)^2} + \frac{1}{(1 - v_0/4)^2} \right]$$

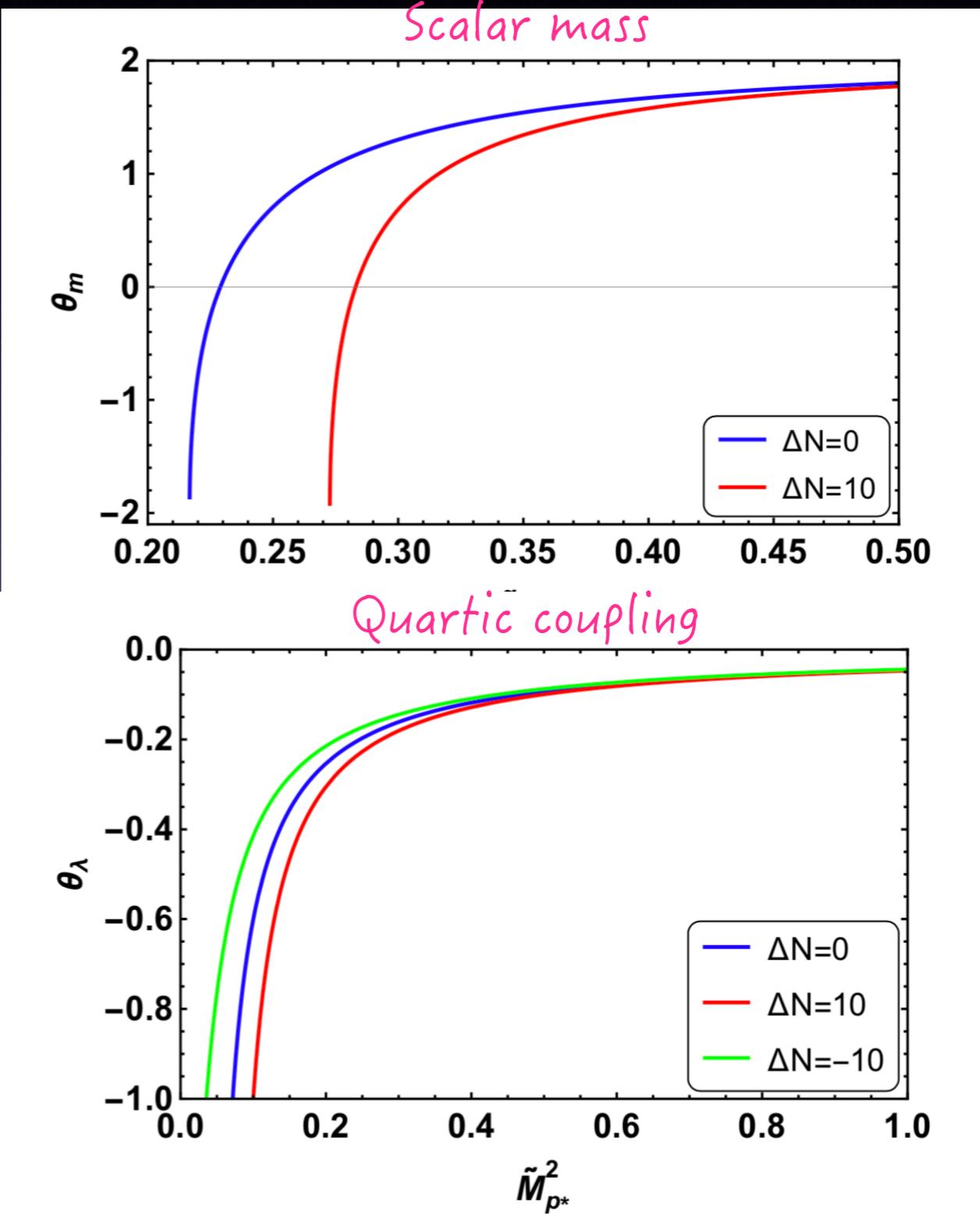
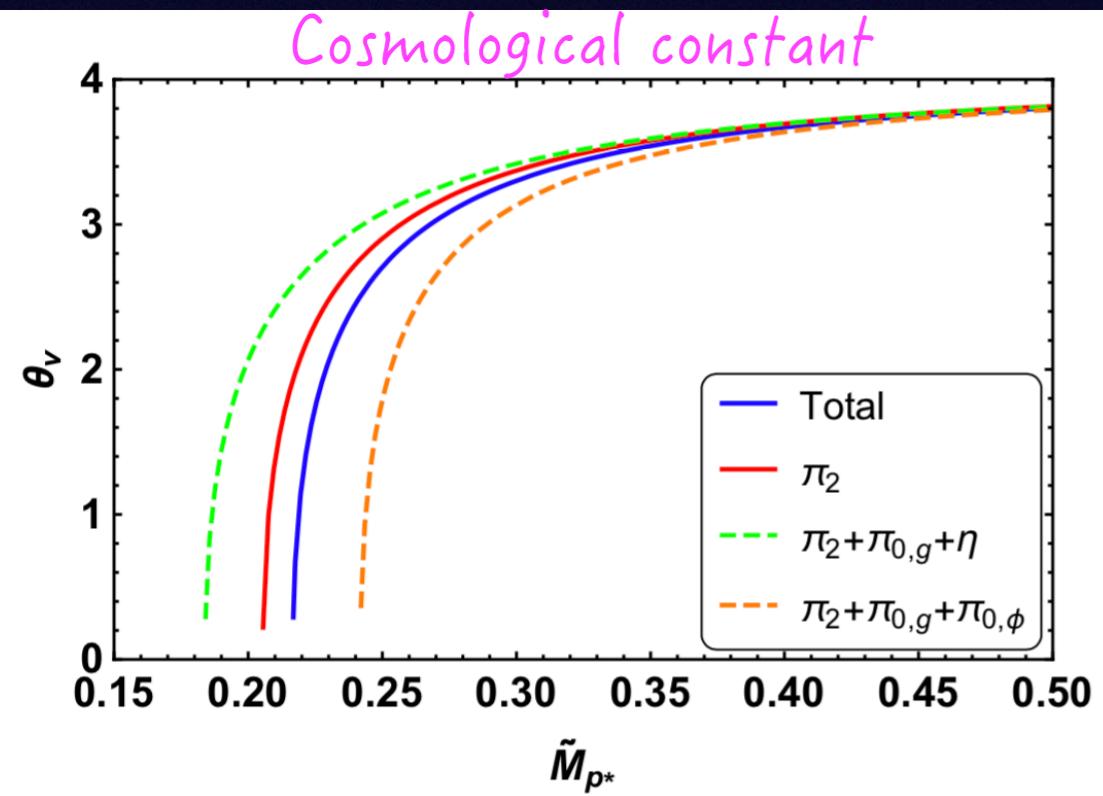
Anomalous dimension

Fixed point of v_{0^*}

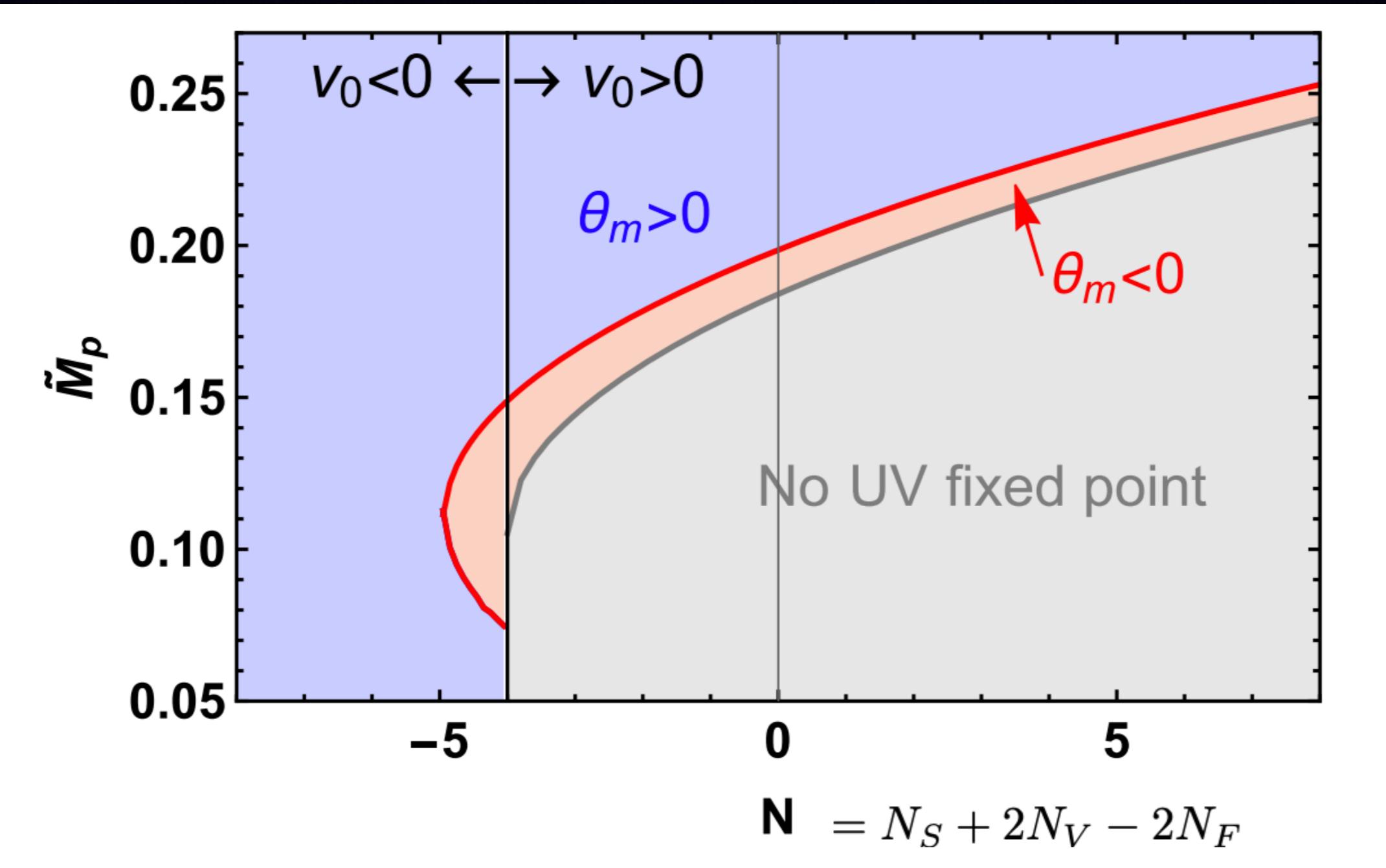


Critical exponents

$N = 4, \Delta N = 0$



N vs. Mp



Summary

- Asymptotically safe gravity induces large anomalous dimension.
- Scaling of Higgs scalar potential above the Planck scale is changed from canonical scaling.
- Phase diagram on (N, M_p) -plane.